

## Shock Tube Studies of the Acetylene and Ethylene Pyrolysis by UV Absorption

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Absorption measurements at 216 and 230 nm of shock heated  $C_2H_2$ ,  $C_2H_2+H_2$ , and  $C_2H_4$  diluted in Ar were made over the temperature range 1800 K to 2600 K at half of atmospheric pressure. Absorptivities were evaluated for  $C_4H_2$ , which is the main product of  $C_2H_2$  pyrolysis, as well as for  $C_2H_2$  and  $C_2H_4$ . The relative value of  $C_4H_2$  absorption was confirmed to decrease with the presence of  $H_2$  in accordance with a computer modeling result using a 25-reaction mechanism. The time profile of  $C_4H_2$  absorption could be modeled well with  $k=10^{12.4} \text{ cm}^3 \text{ mol}^{-1} \text{ s}^{-1}$  for the reaction  $H_2+C_2H=C_2H_2+H$ .

The pyrolysis of  $C_2H_2$  in shock waves has been studied by many workers and has become explainable through computer modeling with a rather definite reaction mechanism.

Recently, two groups studied the  $C_2H_2$  pyrolysis using real-time kinetic absorption spectroscopy and the laser schlieren method. Frank and Just<sup>1)</sup> measured the  $[H]$  profile in the  $C_2H_2$  pyrolysis by atomic resonance absorption spectroscopy (ARAS) in mixtures of  $C_2H_2$  highly diluted in Ar to ppm order and evaluated the initiation reaction rate constant in the low pressure region. Tanzawa and Gardiner<sup>2)</sup> did laser schlieren experiments on the  $C_2H_2$  pyrolysis. They proposed a reaction mechanism, in which the experimental results obtained by TOF mass spectroscopy<sup>3,4)</sup> and single pulse shock tube<sup>5)</sup> as well as their laser schlieren results were analyzed by computer modeling studies. Their later efforts<sup>6)</sup> led them to a final mechanism which accords with almost all of the experimental results from 625 to 3400 K.

The new mechanism of the  $C_2H_2$  pyrolysis so derived is expected to improve the  $C_2H_4$  pyrolysis mechanism, because  $C_2H_4$  decomposes to  $C_2H_2$  and  $H_2$  by a molecular reaction very rapidly in the high temperature region.<sup>7)</sup>

Cundall *et al.*<sup>8)</sup> studied the  $C_2H_2$  and  $C_2H_4$  pyrolysis in shock waves by absorption spectroscopies with Xe-lamp and He-Ne laser. The reaction mechanism used by them was rather primitive and the values of heat of formation for  $C_2H$  and polyacetylene radicals are subject to correction,<sup>1,2,9)</sup> though they could obtain qualitative information about soot formation.

The present paper is concerned with the  $C_2H_2$  and  $C_2H_4$  pyrolysis behind incident shock waves. The reaction features were examined by UV absorption at 216 and 230 nm and were interpreted more quantitatively by the computer modeling with a refined mechanism.<sup>6)</sup>

### Experimental

The shock tube used in this study has been described in detail already.<sup>10)</sup>

The light from a D<sub>2</sub>-lamp (Hamamatsu-TV, L544) or Xe-lamp (Ushio, UXL500D), made parallel with a quartz lens, was measured using a photomultiplier (Hamamatsu-TV, R208) with a monochrometer (Jarrell-Ash, JE25) after passing through two shock tube windows. The output from the photomultiplier was displayed on an oscilloscope (Tektronix, 7603).

The signal from piezo-gauges mounted on the shock tube every 10 cm was detected by a universal counter (Takeda-Riken, TR5002). Shock temperature was calculated from the measured shock velocity assuming full relaxation and no chemical reaction.

Gases used were 99.7% pure  $C_2H_2$  and 99.6% pure  $C_2H_4$  purchased from Takachiho and 99.999% pure Ar from Nippon Sanso. All gases were used without further purification. The test gas was prepared in a 20 l glass bulb manometrically and used after more than 48 h mixing time. The test gas compositions are shown in Table 1.

TABLE 1. MEASURED  $A$  AND MODELED  $[C_4H_2]$  AND  $[C_6H_2]$  AT 2000 K IN EACH TEST GAS MIXTURE (cm and mol units)

No.	Gas composition in Ar	$A$	$[C_4H_2]/10^{-9}$	$[C_6H_2]/10^{-10}$
1	3.2% $C_2H_2$	1.11	15	11
2	2.0% $C_2H_2$ /2.3% $H_2$	0.26	4.3	1.1
3	4.0% $C_2H_4$	0.94	10	2.0

The test section of the shock tube was evacuated to less than  $10^{-4}$  Torr before each run. Neither emission from  $C_2H_2$  or  $C_2H_4$  in Ar mixtures nor absorption by Ar was observed under the present experimental conditions. (1 Torr = 133.322 Pa)

Computer modeling was done using the program mentioned previously.<sup>10)</sup> The flow model used for reflected shock waves was modified and a constant density calculation was adopted.<sup>11)</sup> Thermochemical properties of the species which appear in the reaction mechanism were taken from the JANAF thermochemical table<sup>12)</sup> or other available sources.<sup>9,13,14)</sup>

### Results and Discussion

A sample oscillogram of the UV absorption (216 nm, 3.2%  $C_2H_2$  in Ar, and  $P_2=0.35$  atm) is shown in Fig. 1. We evaluated absorptivities of  $C_2H_2$  and  $C_2H_4$  from the absorption at the shock front for two wavelengths, 216 and 230 nm, by using the following definition:

$$a = \ln(I_0/I)/[C]_0/d,$$

where  $I_0$ =incident light intensity,  $I$ =transmitted intensity at the shock front,  $[C]_0$ = $C_2H_2$  or  $C_2H_4$  concentration at the shock front, and  $d$ =optical path length (7.2 cm). The obtained  $a$  values of  $C_2H_2$  and  $C_2H_4$  are shown in Figs. 2a and 2b.

Acetylene shows an absorption due to the electronic transition  $X \rightarrow A$  in the wavelength region 210–237

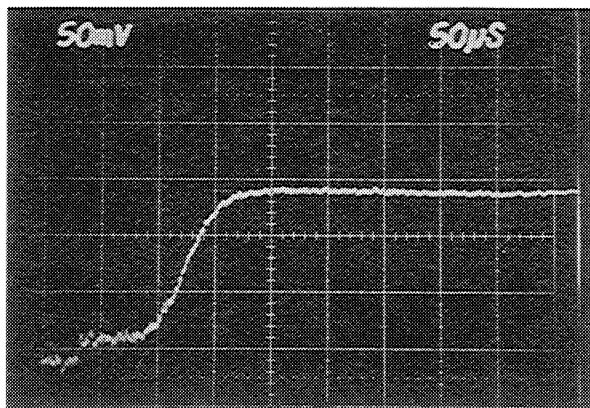


Fig. 1. A representative profile of transmitted D<sub>2</sub>-lamp intensity for 3.2% C<sub>2</sub>H<sub>2</sub> in Ar mixture,  $T_2=2084$  K, and  $P_2=0.35$  atm.

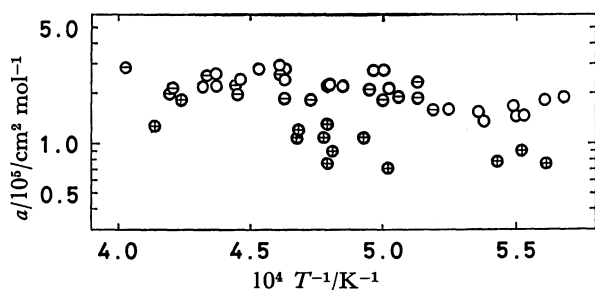


Fig. 2a. Temperature dependence of C<sub>2</sub>H<sub>2</sub> absorptivity at 216 and 230 nm with D<sub>2</sub>-lamp. The symbols are as follows.

○: for 3.2/96.8=C<sub>2</sub>H<sub>2</sub>/Ar at 216 nm, ○⊙: for 2.0/2.2/95.8=C<sub>2</sub>H<sub>2</sub>/H<sub>2</sub>/Ar at 216 nm, ⊕: for 3.2/96.8=C<sub>2</sub>H<sub>2</sub>/Ar at 230 nm.

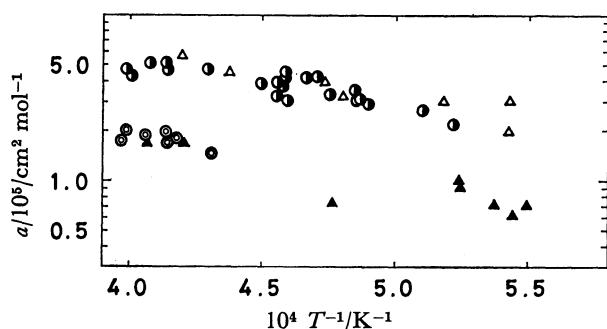


Fig. 2b. Temperature dependence of C<sub>2</sub>H<sub>4</sub> absorptivity at 216 and 230 nm. The symbols are as follows.

●: for 4.0% C<sub>2</sub>H<sub>4</sub> in Ar with D<sub>2</sub>-lamp at 216 nm, ●⊙: for 4.0% C<sub>2</sub>H<sub>4</sub> in Ar with D<sub>2</sub>-lamp at 230 nm, △: for 4.5% C<sub>2</sub>H<sub>4</sub> in Ar with Xe-lamp at 216 nm, ▲: for 4.5% C<sub>2</sub>H<sub>4</sub> in Ar with Xe-lamp at 230 nm. Data taken with Xe-lamp were corrected for higher-order wavelengths.

nm.<sup>15)</sup> In comparison with C<sub>2</sub>H<sub>2</sub>, C<sub>2</sub>H<sub>4</sub> has greater absorptivity in this region, though no absorption which can be attributed to an electronic transition is reported for C<sub>2</sub>H<sub>4</sub>.<sup>15)</sup> Figures 2a and 2b show that the absorptivities of C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>4</sub> at 216 nm are larger than those at 230 nm in the temperature range measured.

This may imply that the absorption of C<sub>2</sub>H<sub>4</sub> at 216 nm is attributed to an extended transition of X→A. Since the wavelength 230 nm is too far from the wavelength region of the X→A transition, the absorption at 230 nm may not be the extended transition. No clear explanations of the C<sub>2</sub>H<sub>4</sub> absorption at 230 nm at these high temperatures can be given only on the basis of the present experimental results.

The  $a$  values of C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>4</sub> increase slightly with increasing temperature, in contrast with the reported  $a$  values of CH<sub>3</sub> which have the opposite behavior.<sup>16,17)</sup> Similar observations are reported recently for 3.39  $\mu$ m He-Ne laser absorption of small hydrocarbons, *i.e.*, the infrared  $a$  values of C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>4</sub> increase with increasing temperature, while those of the alkanes decrease with increasing temperature.<sup>18,19)</sup>

The absorption measurement at 216 nm was employed to study ethane decomposition and CH<sub>3</sub> recombination in shock waves.<sup>16,17)</sup> The results shown above imply that the absorption at 216 nm is not only due to CH<sub>3</sub> but also to C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>4</sub>. For alkane pyrolysis measured at high temperatures, CH<sub>3</sub> is usually formed at an early stage of the reaction and C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>4</sub> gradually accumulate afterwards. The absorption at 216 nm at the early stage of alkane pyrolysis represents the features of CH<sub>3</sub> alone; the later absorption begins to show aggregate features of CH<sub>3</sub>, C<sub>2</sub>H<sub>2</sub>, and C<sub>2</sub>H<sub>4</sub>. Our trial runs for the C<sub>2</sub>H<sub>6</sub> pyrolysis measured by 230 nm absorption, on the other hand, showed negligible absorptivity for CH<sub>3</sub> at the 1st stage of the pyrolysis. This result and the above may explain why we selected 216 and 230 nm absorptions to study the C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>4</sub> pyrolysis. The computer program, which incorporates these absorptivities for C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>4</sub> at 216 and 230 nm and the absorptivity for CH<sub>3</sub> at 216 nm, may solve the alkane pyrolysis measured at the two wavelengths by a procedure similar to that adopted in the C<sub>2</sub>H<sub>6</sub> and C<sub>3</sub>H<sub>8</sub> pyrolysis studies monitored by a 3.39  $\mu$ m He-Ne laser.<sup>14,20)</sup>

The oscillogram of Fig. 1 shows that the absorption increases almost linearly after the absorption at the shock front and finally reaches a high plateau value. The plateau absorption  $A$  was defined by  $A=\ln(I_0/I_t)$ , where  $I_t$  is the transmitted light intensity at the plateau and  $I_0$  the incident light intensity. In Fig. 3, the relative  $A$  value or the value of  $A$  divided by the shock front concentration of [C<sub>2</sub>H<sub>2</sub>]<sub>0</sub> or [C<sub>2</sub>H<sub>4</sub>]<sub>0</sub> are shown *vs.*  $10^4/T$ .

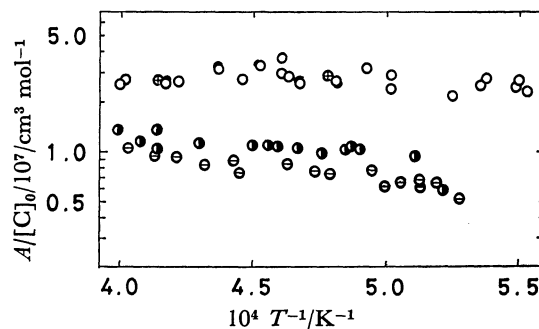


Fig. 3. Temperature dependence of relative steady absorption at 216 and 230 nm. The symbols are as in Figs. 2a and 2b.

In this case, the relative values are free from the choice between the two wavelengths, 216 and 230 nm, but the values show distinct mixture dependence, splitting into two groups as shown in Fig. 3. Searches for the chemical species which are responsible for the plateau absorption were made by the computer modeling, adopting the reaction mechanism shown in Table 2.

TABLE 2. REACTION MECHANISM AND RATE CONSTANTS  
 $k = A \exp(-E/RT)$  (mol, cm, s, and kcal units)

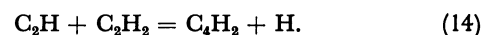
Reaction	log $A$	$E$	Ref. No.
1) $C_2H_4 + M = C_2H_2 + H_2 + M$	17.4	79.3	7
2) $C_2H_4 + M = C_2H_3 + H + M$	17.6	98.2	7
3) $C_2H_4 + C_2H_4 = C_2H_3 + C_2H_5$	14.8	64.0	24
4) $C_2H_4 + H = C_2H_3 + H_2$	15.7	22.9	7
5) $C_2H_3 + C_2H_4 = C_4H_6 + H$	12.0	7.3	25
6) $C_2H_3 + H = C_2H_2 + H_2$	13.0	0.	14
7) $C_2H_5 + M = C_2H_4 + H + M$	15.3	30.1	14
8) $C_2H_2 + M = C_2H + H + M$	16.6	106.5	1
9) $C_2H_2 + C_2H_2 = C_4H_3 + H$	12.3	45.0	6
10) $C_2H_3 + M = C_2H_2 + H + M$	14.9	31.0	25
11) $C_2H_2 + C_2H_3 = C_4H_4 + H$	13.2	25.0	6
12) $C_4H_4 + C_2H = C_6H_3 + C_2H_2$	13.6	0.	6
13) $C_2H + H_2 = C_2H_2 + H$	12.4	0.	This work
14) $C_2H + C_2H = C_4H_2 + H$	13.6	0.	2
15) $C_4H_3 + M = C_4H_2 + H + M$	16.0	60.0	2
16) $C_4H_2 + M = C_4H + H + M$	17.5	80.0	2
17) $C_6H_2 + M = C_6H + H + M$	16.7	80.0	2
18) $C_6H_2 + M = C_6H + H + M$	16.7	80.0	2
19) $H_2 + M = 2H + M$	12.4	92.6	26
+0.5 log $T$			
20) $C_4H_3 + H + M = C_4H_4 + M$	15.0	0.	2
21) $C_2H_2 + C_4H = C_6H_2 + H$	13.6	0.	2
22) $C_4H_2 + C_2H = C_6H_2 + H$	13.6	0.	2
23) $C_2H_2 + C_6H = C_8H_2 + H$	12.0	0.	2
24) $C_6H_2 + C_2H = C_8H_2 + H$	12.0	0.	2
25) $C_4H_2 + C_4H = C_8H_2 + H$	12.0	0.	2

Acetylene radical and polyacetylene radicals can not be the absorbers; if they were, they must have large, temperature-dependent absorptivities. For example, the steady concentration of  $C_2H$  at 2500 K is about 20 times larger than that at 2000 K in 3.2%  $C_2H_2$  in Ar mixtures. This means that the absorptivity of  $C_2H$  decreases by a similar factor between the two temperatures, because the  $A$  values are almost constant over the temperature range measured. We are unaware of any species having such large, temperature-dependent absorptivity.

Table 1 shows the  $A$  values, obtained by a least squares fitting for the  $A/[C]_0$  and  $10^4/T$  relation shown in Fig. 3, and the steady  $[C_4H_2]$  and  $[C_6H_2]$ , obtained by the computer modeling, at 2000 K in different three mixtures. The ratios of the  $A$  values,  $A_m/A_n$ , are not in accord with the  $[C']_m/[C']_n$  values, where subscripts m and n denote the mixture number shown in Table 1 and  $C'$  is  $C_4H_2$  or  $C_6H_2$ . Although these modeled  $[C_4H_2]$  and  $[C_6H_2]$  are subject to the values of heat of formation adopted, no clear-cut explanation with appropriate thermochemical data could be given.

There is an electronic transition designated as  $X \rightarrow B$  over the wavelength range 200 to 265 nm for  $C_4H_2$ .<sup>15)</sup> The reported  $C_4H_2$  profiles in the  $C_2H_2$  pyrolysis measured by TOF mass spectroscopy<sup>4)</sup> are similar to the absorption profile of Fig. 1. These results strongly support the conjecture that the steady absorption is mainly due to  $C_4H_2$ . The absorptivity of  $C_4H_2$  was evaluated as:  $a(\text{cm}^2 \text{mol}^{-1}) = (1.1 \pm 0.3) \times 10^7$  at 2000 K and  $(1.4 \pm 0.2) \times 10^7$  at 2500 K.

The plateau absorption shown in Fig. 1 indicates that the overall reaction of the  $C_2H_2$  and  $C_2H_4$  pyrolysis reaches an equilibrium under the present experimental conditions; this was confirmed also by the computer modeling. The primary reactions of the  $C_2H_2$  pyrolysis in these temperature ranges are:<sup>6)</sup>



Ethynyl,  $C_2H$ , and  $H$  play roles as chain carriers in the reactions. An equilibrium between these species leads to an equilibrium for  $[C_4H_2]$ :

$$[C_4H_2] = K_{-13}K_{14}[C_2H_2]^2/[H_2],$$

where  $K_{-13}$  and  $K_{14}$  are the equilibrium constants of the reactions, (-13) and (14). The above equation is in accordance with the result shown in Fig. 3 that the presence of  $H_2$  inhibits the formation of  $C_4H_2$ , which is one of the precursors to soot.<sup>8)</sup> Relative  $A$  values in the  $C_2H_4/\text{Ar}$  mixture are smaller than those in the  $C_2H_2/\text{Ar}$  mixture. This result also can be explained in the same way. Ethylene decomposes by a molecular reaction fairly rapidly in this temperature range<sup>7)</sup> and forms  $C_2H_2$  and  $H_2$  in equal amounts. Therefore, in the  $C_2H_4/\text{Ar}$  mixture,  $C_4H_2$  is produced by the pyrolysis of  $C_2H_2$  so formed in the presence of  $H_2$ . The fact that the measured relative  $A$  values in the  $C_2H_4/\text{Ar}$  and  $C_2H_2/H_2/\text{Ar}$  mixtures are similar to each other may support the above explanation.

Figure 1 provides us with time parameters whose value can be used as a measure to estimate the overall rate of the  $C_2H_2$  and  $C_2H_4$  pyrolysis. In the conventional manner, the induction time can be defined as: the time between the shock arrival and the foot of the linear absorption rise. If the absorptivity of  $C_4H_2$  is exactly known, the threshold value of  $[C_4H_2]$  at the foot can be estimated and compared with the modeled  $[C_4H_2]$ . But there are still uncertainties in the  $C_4H_2$  absorptivity obtained above. Instead, we defined a parameter,  $t_m$ : time between the shock arrival and the crossing point where the linear absorption rise and the steady absorption are extrapolated to each other. The parameter,  $t_m$ , so defined is independent of the  $C_4H_2$  absorptivity. The modeled  $t_m$  was evaluated from the  $C_4H_2$  profile. First of all, we compared the measured  $t_m$  with the modeled one using the mechanism and the rate constants in Table 2 except for the  $k_{13}$  value in the  $C_2H_2/\text{Ar}$  mixture. As for the  $k_{13}$  value,  $10^{13.54}$  in Ref. 6 was used. Figure 4 shows the relation between  $t_m[C_2H_2]_0$  or  $t_m[C_2H_4]_0$  and  $10^4/T$  for the measured and modeled results. As is shown by the broken line, the modeled  $t_m$  was found to be smaller than that measured by a similar factor over the temperature range measured. To confirm the correctness of the interpretations for the

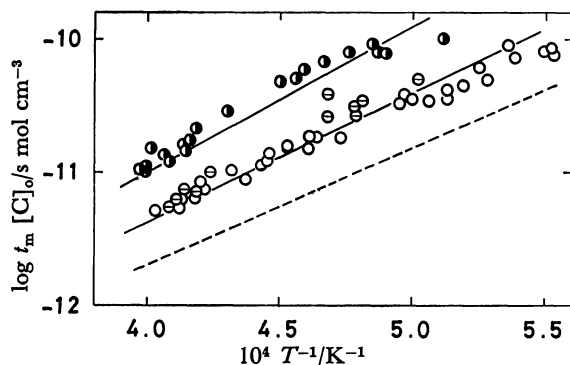


Fig. 4.  $t_m[C_2H_2]_0$  or  $t_m[C_2H_4]_0$  vs.  $10^4/T$ . The symbols are the same as in Figs. 2a and 2b. The solid lines show modeled value using Table 2 mechanism and rate constants. The upper one is for  $C_2H_4/Ar$  mixture and the lower one is for both of  $C_2H_2/Ar$  and  $C_2H_2/H_2/Ar$  mixtures. The broken line shows modeled value for 3.2%  $C_2H_2$  in Ar mixture using Table 2 mechanism and rate constants except for the  $k_{13}$  value:  $k_{13}=10^{13.63}$  was used.<sup>6)</sup>

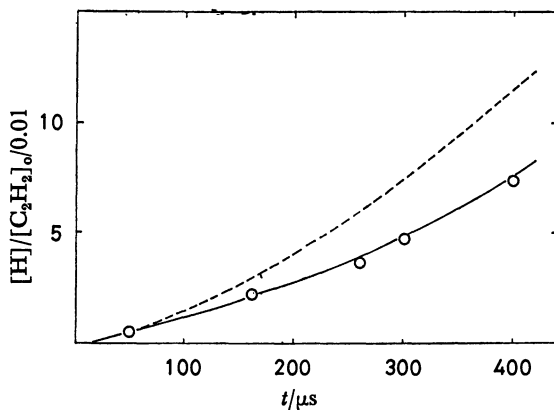


Fig. 5. Comparison measured  $[H]$  profile<sup>1)</sup> with modeled ones. The circles are data points of measured  $[H]$  profile in reflected shock wave for 20 ppm in Ar mixture at  $T_s=2450$  K and  $[Ar]_s=8.2 \times 10^{-6}$ . The solid and broken lines are modeled values as are in Fig. 4.

present experiment and modeling, another  $C_2H_2$  pyrolysis result was adopted for the modeling.

Gas sampling problems due to boundary layer growth are present in the results obtained by the single pulse shock tube technique.<sup>21)</sup> Since TOF mass spectroscopy may also have the same problems, we selected the results obtained by ARAS<sup>1)</sup> for the modeling. Figure 5 shows both the measured and the modeled  $[H]$  profiles. It is clear that the modeled  $[H]$  increasing rate shown by the broken line is larger than the measured one, which is in accordance with the result obtained from Fig. 4, i.e., the modeled overall reaction rate is faster than the measured one.

The first mechanism proposed by Tanzawa and Gardiner,<sup>2)</sup> which was derived to explain the high temperature pyrolysis of  $C_2H_2$  in shock waves, could model the results of TOF mass spectroscopy<sup>3,4)</sup> and single pulse shock tube<sup>5)</sup> fairly well, whereas the new mechanism in Ref. 6 gave rather poor modeling results. The new mechanism in Ref. 6 seems to have an improper

TABLE 3. SENSITIVITY SPECTRUM WITH PS VALUES<sup>a)</sup>

Reaction	$t_m$ <sup>b)</sup>	$[H]_{100 \mu s}$ <sup>c)</sup>	$[H]_{400 \mu s}$ <sup>c)</sup>
$C_2H_2 + M = C_2H + H + M$	-8/-5 <sup>d)</sup>	87/97	54/87
$C_2H_2 + C_2H_2 = C_4H_3 + H$	-8/-5	—	—
$C_2H + H_2 = C_2H_2 + H$	-29/-42	30/47	29/57
$C_2H_2 + C_2H = C_4H_2 + H$	-3/-18	11/3	25/15
$C_4H_2 + M = C_4H + H + M$	-21/-27	—	—

a)  $PS = \log(\text{parameter}'/\text{parameter})/\log(\text{multiplier}) \times 100$ . Details are shown in Ref. 22. b) For 3.2%  $C_2H_2$  in Ar mixture,  $P_2=0.37$  atm, and  $T_2=2200$  K. c) For 50 ppm  $C_2H_2$  in Ar mixture,  $P_5=1.77$  atm, and  $T_5=2420$  K. d) The first entry is obtained by multiplying each rate constant by 5; the second entry, by 0.2.

value for some rate constant. To find which reaction is responsible to the disagreement shown in Figs. 4 and 5, a sensitivity calculation with PS values<sup>22)</sup> was tried for our experiment and for the ARAS measurement.<sup>1)</sup>

Table 3 shows the reactions having large PS values and the obtained PS values for the parameters. The reaction which has the largest PS values is different in our experiment and in the ARAS experiment. The effect of secondary reactions on the  $[H]$  profile is suppressed in the latter case, in which highly diluted test gas mixtures were used, and the measured and modeled  $[H]$  profiles reflect the initiation reaction,  $C_2H_2 + M = C_2H + H + M$ , especially at the early stages of the reaction. The PS values of reaction 14 becomes large as the reaction proceeds, in contrast to those of the initiation reaction, while reaction 13 has an almost constant PS values. The time dependent variations of PS values imply that the chain carrier concentrations such as H and  $C_2H$  increase and the chain reactions dominate. These effects of the chain reactions on the parameters adopted are very distinct in our experiment.

Reaction 16,  $C_4H_2 + M = C_4H + H + M$ , has the second largest PS values for  $t_m$  in our experiment. Frank and Just<sup>1)</sup> measured the  $k_{16}$  value in the high pressure limit directly by monitoring the  $[H]$  profile produced by the  $C_4H_2$  pyrolysis using the ARAS. Although it is possible to derive the  $k_{16}$  value suitable for the modeling of our experiment from the reported  $k_{16}$  value by making a reduced fall-off curve,<sup>23)</sup> we left the  $k_{16}$  value used in this study at the value determined indirectly by Tanzawa and Gardiner,<sup>2)</sup> through the computer modeling for the laser schlieren experiment because the experimental conditions selected by them are similar to the present ones except for the temperature range.

Since the  $C_2H_2$  pyrolysis is almost equilibrated at the reaction time when the parameter  $t_m$  is measured, it is to be expected that reaction 13 has the largest PS values. So only the  $k_{13}$  value was changed so as to fit the modeled  $t_m$  to the measured over the temperature range measured. An excellent agreement between the experiment and the modeling for the  $[H]$  profiles, as well as for the  $t_m$  values, was obtained by adopting the following  $k_{13}$  value:

$$k_{13} = 10^{12.4} (\text{cm}^3 \text{mol}^{-1} \text{s}^{-1}).$$

This is shown in Figs. 4 and 5. In our computer

program,<sup>10)</sup> reverse rate constant of a given elementary reaction is calculated using an equilibrium constant; an assembly of first order differential equations of both forward and reverse reactions is numerically integrated. Then, when either  $k_{13}$  or  $k_{-13}$ , where  $K_{-13}=k_{-13}/k_{13}$ , is incorporated into the reaction mechanism with reasonable thermochemical data, the same modeling results can be obtained. Tanzawa and Gardiner<sup>2)</sup> proposed a non-Arrhenius temperature-dependent  $k_{-13}$  value in the first mechanism. The  $k_{13}$  value so evaluated from the  $k_{-13}$  value was found to have exactly the same value at 2200 K as the present  $k_{13}$  value. Then, it is not surprising that the  $k_{13}$  value obtained in this study is about 10 times smaller than that in Ref. 6.

We note that the modeling result of the  $C_2H_2$  and  $C_2H_4$  pyrolysis is affected by the  $\Delta H_f^\circ$  value adopted for acetylene and polyacetylene radicals. Two recent papers about the  $C_2H_2$  pyrolysis study<sup>1,2)</sup> demonstrated the correctness of the new  $\Delta H_f^\circ$  for  $C_2H$ .<sup>9)</sup> As has been shown already, this new value of  $\Delta H_f^\circ$  for  $C_2H$  will be a help to improve the  $\Delta H_f^\circ$  values for polyacetylene radicals using the additivity rule.<sup>27)</sup> The  $\Delta H_f^\circ$  value for  $C_4H$  so derived is 180 kcal/mol,<sup>1)</sup> which is 26 kcal/mol larger than the old value.<sup>13)</sup> If the old value for  $C_4H$  is adopted, unusual profiles of  $C_4H_2$  and  $C_4H$  can be obtained, e.g., steady  $[C_4H]$  is larger than steady  $[C_4H_2]$ .

We have little to say about the  $C_2H_4$  pyrolysis during the early reaction period on the basis of the present experimental results. The computer modeling for the  $C_2H_4/Ar$  mixture showed that almost all of  $C_2H_4$  is pyrolyzed to  $C_2H_2$  and  $H_2$  already at the reaction time when the linear absorption starts.

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† 1 kcal=4.184 kJ.